

TPV

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Blatt 3

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Exercise 6 *Free electromagnetic field as a harmonic oscillator*

The free electromagnetic field in the Coulomb gauge ($\nabla \cdot \mathbf{A}(\mathbf{r}, t) = 0$) is fully described by a vector potential $\mathbf{A}(\mathbf{r}, t)$ which can be decomposed in the following form:

$$\mathbf{A}(\mathbf{r}, t) = \sum_{\lambda} q_{\lambda}(t) \mathbf{f}_{\lambda}(\mathbf{r}). \quad (1)$$

For a box of volume V and length L , the functions $\mathbf{f}_{\lambda}(\mathbf{r})$ are given by

$$\mathbf{f}_{\lambda}(\mathbf{r}) = \sqrt{\frac{8\pi c^2}{L^3}} \hat{\epsilon}_{\lambda} \begin{cases} \cos(\mathbf{k}_{\lambda} \cdot \mathbf{r}), \\ \sin(\mathbf{k}_{\lambda} \cdot \mathbf{r}), \end{cases} \quad \mathbf{k}_{\lambda} = \frac{2\pi}{L} (n_x, n_y, n_z) \quad (2)$$

with $n_x, n_y \in \mathbb{Z}$, $n_z \in \mathbb{N}^*$ and possible polarizations $\hat{\epsilon}_{\lambda} = \hat{\epsilon}_{\lambda}^{(1)}, \hat{\epsilon}_{\lambda}^{(2)} \perp \mathbf{k}_{\lambda}$. The $\mathbf{f}_{\lambda}(\mathbf{r})$ form a complete orthonormal basis satisfying the orthogonality relation

$$\int \mathbf{f}_{\lambda'}(\mathbf{r}) \mathbf{f}_{\lambda}(\mathbf{r}) d\mathbf{x} = \delta_{\lambda'\lambda} \mathcal{N}, \quad (3)$$

\mathcal{N} being some normalization constant. For the given $\mathbf{f}_{\lambda}(\mathbf{r})$ the Coulomb-gauge condition follows directly. If the vector potential satisfies the periodic boundary conditions the functions $\mathbf{f}_{\lambda}(\mathbf{r})$ will be solutions to the following equation:

$$\nabla^2 \mathbf{f}_{\lambda}(\mathbf{r}) + \frac{\omega_{\lambda}^2}{c^2} \mathbf{f}_{\lambda}(\mathbf{r}) = 0 \quad (4)$$

- a) Write an Hamiltonian density for the free electromagnetic field $\mathcal{H} = \frac{1}{8\pi} (|\mathbf{E}|^2 + |\mathbf{B}|^2)$ in terms of $q_{\lambda}(t)$ and $\mathbf{f}_{\lambda}(\vec{r})$ functions.

(1 Point)

- b) Prove that the Hamiltonian of the electromagnetic field $H = \int_V \mathcal{H} d\mathbf{r}$ can be written as the Hamiltonian of a collection of harmonic oscillators:

$$H = \sum_{\lambda} \left(\frac{1}{2} p_{\lambda}^2 + \frac{\omega_{\lambda}^2}{2} q_{\lambda}^2 \right) \quad (5)$$

Hint 1 : To derive the Hamiltonian use the relation:

$$\nabla \cdot (\mathbf{A} \times (\nabla \times \mathbf{B})) = (\nabla \times \mathbf{A}) \cdot (\nabla \times \mathbf{B}) - \mathbf{A} \cdot (\nabla \times (\nabla \times \mathbf{B})). \quad (6)$$

Hint 2 : Use the Ostrogradsky-Gauss theorem (divergence theorem) and the periodic boundary conditions: $\mathbf{f}_{\lambda}(x, y, z) = \mathbf{f}_{\lambda}(x + L, y, z) = \mathbf{f}_{\lambda}(x, y + L, z) = \mathbf{f}_{\lambda}(x, y, z + L)$.

(3 Points)

Exercise 7 *The energy-momentum tensor*

a) Show that the energy-momentum tensor density defined by:

$$\mathcal{T}_{\mu\nu} = -\frac{\partial\phi}{\partial x_\nu} \frac{\partial\mathcal{L}}{\partial(\partial\phi/\partial x_\mu)} + \mathcal{L}\delta_{\mu\nu} \quad (7)$$

satisfies the continuity equation:

$$\frac{\partial\mathcal{T}_{\mu\nu}}{\partial x_\mu} = 0. \quad (8)$$

when the Euler-Lagrangian equation for ϕ is assumed. *(1 Point)*

b) Show that each component of the four-vector

$$P_\mu(t) = -i \int \mathcal{T}_{4\mu} d^3\mathbf{x} \quad (9)$$

is constant in time if ϕ vanishes sufficiently rapidly at infinity. (The integration is over three dimensional space at a given instant t .) *(1 Point)*

c) Derive the Hamiltonian density \mathcal{H} and show that for the real scalar field it holds:

$$\mathcal{H} = -\mathcal{T}_{44}. \quad (10)$$

(1 Point)