

11. Thermodynamic quantities of paramagnets

We consider a system of N non-interacting spin-1/2 particles which are in a static magnetic field $\vec{B} = B_0 \vec{e}_z$. The Hamiltonian operator is then given by

$$\hat{H} = -B_0 \hat{M}^z = \mu_B B_0 \sum_{j=1}^N \hat{\sigma}_j^z \quad (1)$$

where \hat{M}^z is the magnetisation operator in z -direction and μ_B is the Bohr magneton. Furthermore, let $\hat{\rho}_N$ be the density matrix of the entire system of N spins.

- (a) Show that at equilibrium and in a homogeneous system the following statement is true

$$U = \langle \hat{H} \rangle = N \text{Tr} \{ \hat{\rho}_1 \hat{H}_1 \} \quad (2)$$

with $\hat{H}_1 = \mu_B B_0 \hat{\sigma}_1^z$ and $\hat{\rho}_1 = \text{Tr}_{N-1} \{ \hat{\rho}_N \}$, and Tr being the trace over the $N - 1$ other spins.

(1 point)

- (b) Assume a Boltzmann-Gibbs distribution of spins. The density matrix of a single spin is then given by

$$\hat{\rho}_1 = \frac{e^{\beta \epsilon_0} |+\rangle \langle +| + e^{-\beta \epsilon_0} |-\rangle \langle -|}{\mathcal{Z}_1}, \quad (3)$$

with the normalisation constant $\mathcal{Z}_1 = e^{\beta \epsilon_0} + e^{-\beta \epsilon_0}$ and $\epsilon_0 = \mu_B B_0$. Show that the total energy U is given by

$$U = -n \mu_B B_0 \frac{e^{2\beta \mu_B B} - 1}{e^{2\beta \mu_B B} + 1} \quad (4)$$

with $\beta = 1/k_B T$ by explicitly determining the expected value of the Hamiltonian $\langle \hat{H} \rangle$. Determine its limits at $T \rightarrow 0$ and $T \ll \mu_B B_0 / k_B$. Show further that the expectation value of the magnetisation in z -direction is given by

$$\langle \hat{M}^z \rangle = N \mu_B \tanh \left(\frac{\mu_B B_0}{k_B T} \right) \quad (5)$$

(1 point)

- (c) Use the explicit form of the probabilities $p_{\pm} = e^{\mp\beta\epsilon_0}/\mathcal{Z}_1$ to derive the following formula for the entropy S of the system.

$$S = -Nk_B[p_+ \ln p_+ + p_- \ln p_-]. \quad (6)$$

(1 point)

- (d) Use the formulas of entropy S or total energy U to derive a formula for the specific heat C and show that

$$C = \left. \frac{\partial U}{\partial T} \right|_{B_0} = T \left. \frac{\partial S}{\partial T} \right|_{B_0} = k_B N (\beta \mu_B B_0)^2 \operatorname{sech}^2(\beta \mu_B B_0). \quad (7)$$

How does the specific heat behave for temperatures in the limit cases $T \ll \mu_B B_0/k_B$ and $T \gg \mu_B B_0/k_B$. Show that the specific heat $C(T)$ reaches its maximum approximately at the Curie temperature $\theta = \mu_B B_0/k_B$.

Hint: You may use that the equation $x \tanh x = 1$ für $x > 0$ has only one solution at $x \approx 1.2$.

(1 point)

12. Classical approximation of the quantum gas

A gas of quantum particles of mass m at a density n and temperature T can be approximated by a classical gas if the mean distance d between the particles, defined as

$$d = \frac{1}{n^{1/3}}, \quad (8)$$

is much larger than the thermal de Broglie wavelength, namely $d \gg \lambda_{\text{th}}$. For a gas of $N = 10^{19}$ particles confined in a volume $V = 1 \text{ cm}^3$, determine the corresponding interparticle distance d and compare it with the de Broglie wavelength of a hydrogen molecules at $T \approx 290 \text{ K}$: $\lambda_{\text{th}} \approx 0.72 \times 10^{-8} \text{ cm}$. Check whether the condition to treat the gas as a classical system is fulfilled or not.

(1 point)

13. Density operator

Consider the Hilbert space \mathfrak{H} onto which the density operator $\hat{\rho} : \mathfrak{H} \rightarrow \mathfrak{H}$ is defined such that it can be written as

$$\hat{\rho} = \sum_{\lambda} p_{\lambda} |\psi_{\lambda}\rangle \langle \psi_{\lambda}|, \quad (9)$$

where for all value of λ the vector $|\psi_\lambda\rangle$ are normalized states of the Hilbert space \mathfrak{H} . Furthermore, the coefficients $p_\lambda \geq 0$ are such that $\sum_\lambda p_\lambda = 1$.

(a) Show that the expectation value of some hermitian operator \hat{A} is given by

$$\langle \hat{A} \rangle = \text{Tr}\{\hat{\rho}\hat{A}\}, \quad (10)$$

where the trace is performed over a basis $\{|n\rangle\}$ of the Hilbert space \mathfrak{H} .

(1 point)

(b) Show that the density operator is Hermitian, namely that $\rho^\dagger = \rho$.

(1 point)

(c) Show that the density operator is definite positive, or in other words that

$$\langle \psi | \hat{\rho} | \psi \rangle \geq 0, \quad (11)$$

for any arbitrary state $|\psi\rangle$ of the Hilbert space.

(1 point)

(d) Show that the density operator is normalized: $\text{Tr}\{\hat{\rho}\} = 1$.

(1 point)