

Exercises for Theoretical physics IV

WiSe 2022/23

Sheet 5

08.11.2022

Exercise 11 *Poisson distribution und Stirling formula*

We consider N independent coin tosses with a flipped coin. Heads occur with probability p , tails with probability $1 - p$. The probability of finding n times heads after N coin tosses is given by the binomial distribution

$$W_N(n) = \binom{N}{n} p^n (1-p)^{N-n} \quad (1)$$

Here we focus on the limiting case of small probabilities of occurrence $p \ll 1$ (condition A) and large numbers of trials $n \ll N$ (condition B). We now want to find an approximation of the binomial distributions in this limiting case. Proceed as follows:

a) Derive the Stirling formula

$$n! \approx \sqrt{2\pi n} \left(\frac{n}{e}\right)^n, \text{ für } n \gg 1 \quad (2)$$

as the approximation of the factorial for large numbers n . Use the gamma function as a continuous continuation of the factorial.

$$n! = \Gamma(n+1) = \int_0^{+\infty} dx e^{n \ln(x)-x} \quad (3)$$

and substitute $y = x/n$. Approximate the remaining integral using the saddle point approximation

$$\int_a^b dx e^{nf(x)} \approx \sqrt{\frac{2\pi}{n|f''(x_0)|}} e^{nf(x_0)}, \text{ für } n \gg 1, \quad (4)$$

for f twice continuously differentiable, arbitrary endpoints $a < b$, and x_0 the global maximum of f .

(1 point)

b) Show that $(1-p)^{N-n} \approx e^{-Np}$.

(1 point)

c) Show that $N!/(N-n)! \approx N^n$

(1 point)

d) Using the previous part of the exercise, derive the Poisson distribution

$$W_N(n) = \frac{\lambda^n e^{-\lambda}}{n!} \quad (5)$$

as a limiting case of the binomial distribution under conditions A and B, i.e. for a large number of trials and small probabilities of occurrence p , where $\lambda = Np$. Calculate the expectation value and variance of the distribution.

(1 point)

Exercise 12 Central limit theorem

The aim of this exercise is to derive the central limit theorem of probability theory. For this purpose we consider N continuous random variables that can take the values $x_i \in \mathbb{R}$ ($i \in \{1, \dots, N\}$), where $\mathbf{x}^T = (x_1, \dots, x_N) \in \mathbb{R}^N$. The probability that the random variables assume values in the "hypercube" $[x_i, x_i + dx_i]$ is given by $w(\mathbf{x}) d^N x = w(x_1, \dots, x_N) d^N x$ given. Furthermore, let $f(\mathbf{x}) \in \mathbb{R}$ be a known function of the random variable \mathbf{x} . Then the expected value (mean value) of this function with respect to the random variables is

$$\langle f(\mathbf{x}) \rangle = \int_{\mathbb{R}^N} f(\mathbf{x}) w(\mathbf{x}) d^N x.$$

Specifically referred to as

$$\langle x_i^n \rangle = \int_{\mathbb{R}^N} x_i^n w(\mathbf{x}) d^N x$$

the n -th moment ($n \in \mathbb{N}$) of the random variable x_i , with $\langle x_i \rangle$ as the corresponding mean and $(\Delta x_i)^2 = \langle (x_i - \langle x_i \rangle)^2 \rangle$ as variance (fluctuation square). Finally we define the characteristic function of the probability density $w(\mathbf{x})$ according to

$$\chi(\mathbf{k}) = \langle e^{-i\mathbf{k} \cdot \mathbf{x}} \rangle = \int_{\mathbb{R}^N} e^{-i\mathbf{k} \cdot \mathbf{x}} w(\mathbf{x}) d^N x,$$

so that conversely we get the probability density $w(\mathbf{x})$ as a Fourier transform of the characteristic function $\chi(\mathbf{k})$ over

$$w(\mathbf{x}) = \frac{1}{(2\pi)^N} \int_{\mathbb{R}^N} e^{i\mathbf{k} \cdot \mathbf{x}} \chi(\mathbf{k}) d^N k.$$

- a) Based on the random variables \mathbf{x} , the function $\bar{f} = f(\mathbf{x}) \in \mathbb{R}$ can also be interpreted as a random variable. Show with the help of the characteristic function that its probability density $\bar{w}(\bar{f})$ results from

$$\bar{w}(\bar{f}) = \langle \delta(\bar{f} - f(\mathbf{x})) \rangle = \int_{\mathbb{R}^N} \delta(\bar{f} - f(\mathbf{x})) w(\mathbf{x}) d^N x.$$

(2 points)

- b) For the sake of simplicity, in this question we restrict ourselves to the one-dimensional case with $x \in \mathbb{R}$ as the only random variable and $w(x)$ as the associated probability density. Show that the characteristic function $\chi(k)$ with $k \in \mathbb{R}$ can be written as a power series provided that all moments $\langle x^n \rangle$ exist:

$$\chi(k) = \sum_{n=0}^{\infty} \frac{(-ik)^n}{n!} \langle x^n \rangle$$

In contrast, the so-called cumulants C_n ($n \in \mathbb{N}$) of the probability density $w(x)$ are defined via the power series expansion of the logarithm of the characteristic function:

$$\ln(\chi(k)) = \sum_{n=1}^{\infty} \frac{(-ik)^n}{n!} C_n.$$

Show that the first three cumulants can be determined from the first three moments of the probability density $w(x)$

$$\begin{aligned} C_1 &= \langle x \rangle, \\ C_2 &= \langle x^2 \rangle - \langle x \rangle^2 = (\Delta x)^2, \\ C_3 &= \langle x^3 \rangle - 3\langle x^2 \rangle \langle x \rangle + 2\langle x \rangle^3 \end{aligned}$$

(2 points)

- c) In the following we consider again the general case N of different random variables and define the so-called covariance matrix whose diagonal elements correspond to the variances $(\Delta x_i)^2$

$$V_{ik} = \langle (x_i - \langle x_i \rangle)(x_k - \langle x_k \rangle) \rangle,$$

The off-diagonal elements of the covariance matrix are a measure of the extent to which the fluctuations of the random variables x_i and x_k are correlated around the mean value. Show that for a probability density of the form $w(\mathbf{x}) = w_1(x_1) \tilde{w}(x_2, \dots, x_N)$ all elements V_{1k} ($k \in \{2, \dots, N\}$) vanish.

(1 point)

- d) We now come to the proof of the central limit theorem. For this we assume that the individual random variables x_i have the different probability densities $w_i(x_i)$ and are completely uncorrelated, i.e. the total probability density results from $w(\mathbf{x}) = w_1(x_1) w_2(x_2) \dots w_N(x_N)$.

The central limit theorem states that the probability density $\bar{w}(\bar{y}) = \langle \delta(\bar{y} - y(\mathbf{x})) \rangle$ of the mean value $y(\mathbf{x}) = \frac{1}{N} \sum_{i=1}^N x_i$ of the various random variables for large N converges to the Gaussian function

$$\bar{w}(\bar{y}) \approx \frac{1}{\sqrt{2\pi(\Delta\bar{y})^2}} e^{-\frac{(\bar{y}-\langle\bar{y}\rangle)^2}{2(\Delta\bar{y})^2}} \quad (N \gg 1) \quad (6)$$

The mean value and the variance of the Gaussian distribution are given by $\langle \bar{y} \rangle = \langle y(\mathbf{x}) \rangle$ and $(\Delta\bar{y})^2 = \frac{1}{N}(\Delta y)^2$, where

$$\langle y(\mathbf{x}) \rangle = \frac{1}{N} \sum_{i=1}^N \langle x_i \rangle \quad \text{und} \quad (\Delta y)^2 = \frac{1}{N} \sum_{i=1}^N [\langle x_i^2 \rangle - \langle x_i \rangle^2]$$

That means, in the limit case $N \rightarrow \infty$ the value of the random variable becomes more or less sharp, since it applies to their relative fluctuation

$$\frac{\Delta\bar{y}}{\langle \bar{y} \rangle} = \frac{1}{\sqrt{N}} \frac{\Delta y}{\langle y(\mathbf{x}) \rangle} \rightarrow 0 \quad \text{für} \quad N \gg 1$$

In particular, the Gaussian function (6) for $N \gg 1$ tends towards the delta function $\delta(\bar{y} - \langle \bar{y} \rangle)$.

Prove the central limit theorem by deriving Eq. (6) under the assumption that all moments $\langle x_i^n \rangle$ exist for the individual probability densities $w_i(x_i)$. Take them into account in your calculation of all terms in the highest two orders of $1/N$ and neglect all higher orders.

(3 points)